## How can a quantum computers solve PDEs?

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### Outline

Introduction

Motivation

Fundamental Question

Fluid Mechanics and Physics

2 Theory of Physics-Informed Neural Networks (PINNs).

Deep Neural Networks

Automatic Differentiation (AD)

Approximation for PINNs

3 PINNs for Solving PDEs.

PINNs formulation for PDEs.

4 Quantum CFD

Quantum Computing

HHL Algorithm

Quantum Variational Approach

Quantum Annealing

6 Hypothesis

Project

Quantum Neural Network

6 Project Challenges



### Significance

- Many physical phenomena are fundamentally described by partial differential equations:
  - Quantum mechanics

Source: https://www.youtube.com/watch?v=3YmeajE-TT8 a

$$\mathcal{F}(\mathbf{u}(\mathbf{z}); \mathbf{q}) = \mathbf{f}(\mathbf{z}) \quad \mathbf{z} \in \Omega,$$

$$\mathcal{B}(\mathbf{u}(\mathbf{z})) = \mathbf{g}(\mathbf{z}) \quad \mathbf{z} \in \partial \Omega, \quad \mathbf{z} = [\mathbf{x}, t] \quad (1)$$

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  - Environmental science

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# **Motivation**: Where are the intensive partial differential equations (PDE)?

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  - Fluid mechanics (Navier-Stokes Equations)

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  - Astrophysics
  - Biology, to cite a few.

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# Turbulence from theory to applications

### Turbulence Fundamental understanding

Turbulent Flow Simulation at the Exascale: Opportunities and Challenges Workshop

Sponsored by
U.S. Department of Energy
Office of Science
Office of Advanced Scientific Computing Research

### Program Committee

Michael Sprague (chair, National Renewable Energy Laboratory)
Staniskie Boldyrev (University of Wisconsin-Madison)
Paul Flischer (Argonne National Laboratory and University of Illinois at Urbana-Champaign)
Ray Grout (National Renewable Energy Laboratory)
William I, Gustafson Jr. (Pacific Northwest National Laboratory)
Robert Moser (University of Texas at Austin)

DOE/ASCR Points of Contact Michael Martin and Robinson Pino

### Abstract

This report details the findings and recommendations from the Terbeleat Five Simulation of the Ennoual Opportunities and Challenger Workshop with two keeled Nagas, 4.5, 2015, and the Research (ASCR). The workshop objectives were to define and describe the challenges and opportunities that comparing at the excased will thing to truthenth few simulations in applied science and into the comparing at the excased will thing to truthenth few simulations in applied science and into the comparing the excased will be supported by the comparing the exception of the science and engineering portfolios, including combustion, plasma physics, nuclear research physics, with every, and attempts are science. The high-pair branged requirits product to the product of the comparing the except and the comparing the except of the except

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<sup>a</sup> F.X.Trias, A. Gorobets, and A. Oliva. "Turbulent flow around a square cylinder Computers Fluids. 123:87–98. 2015.



# Intensive PDE application: Climate Science

 Global atmospheric models for large gap in spatial scales between convection (100km) and global motions  $(1 \times 10^4)$  km;

GEOPHYSICAL RESEARCH LETTERS, VOL. 40, 4922-4926, doi:10.1002/grl.50944, 2013

### Deep moist atmospheric convection in a subkilometer global simulation

Yoshiaki Miyamoto, 1 Yoshiyuki Kajikawa, 1 Ryuji Yoshida, 1 Tsuyoshi Yamaura, 1 Hisashi Yashiro. 1,2 and Hirofumi Tomita 1,2

Received 19 July 2013: revised 6 September 2013: accepted 6 September 2013: published 20 September 2013.

[1] Deep moist atmospheric convection is a key element of called cumulus parameterization [Arakawa, 2004 and referthe weather and climate system for transporting mass. momentum, and thermal energy. It has been challenging to simulate convection realistically in global atmospheric models because of the large gap in spatial scales between convection (100 km) and global motions (104 km). We conducted the first ever subkilometer global simulation and described the features of convection. Through a series of grid-refinement resolution testing, we found that an essential change for convection statistics occurred around 2 km grid spacing. The convection structure, number of convective cells, and distance to the nearest convective cell dramatically changed at this resolution. The convection core was resolved using multiple grids in simulations with grid spacings less than 2.0 km. Citation: Miyamoto, Y., Y. Kajikawa, R. Yoshida, T. Yamaura, H. Yashiro, and H. Tomita (2013). Deep moist atmospheric convection in a subkilometer global simulation, Geophys. Res. Lett., 40, 4922-4926, doi:10.1002/grl.50944.

### 1. Introduction

[2] Deep moist convection plays an important role in transporting energy in the troposphere. The temporal and spatial scales of the convection are 30 min-1 h and 1-10 km, respectively [Emanuel, 1994; Houze, 1994]. Convection also contributes to energy transport from equatorial regions to polar

ences therein). Recent advances in computer power and development of a new type of model to solve fluid motions on a sphere have made it possible to conduct global simulations without cumulus parameterizations [Tomita and Satoh, 2004; Satoh et al., 2008]. Previous studies demonstrated that cloudy atmospheric disturbances can be simulated accurately using such a global model [Miura et al., 2007; Sato et al., 2009; Fudeyasu et al., 2010a; 2010b; Nasuno and Satoh, 2011; Kinter et al., 20131

[4] However, their grid spacings (several kilometers) are still coarser than or comparable to the observed convection scale. It is essentially desired to conduct the simulations with resolution higher than the observed convection scale. Furthermore, although the resolution of several kilometers will be widely used for global simulations in the near future, the resolution dependencies of convections simulated in global models are not vet clear. Previous numerical studies using a limited numerical domain, but without parameterization, have demonstrated the resolution dependence of convection features [Weisman et al., 1997; Bryan and Fritsch, 2002; Petch et al., 2002; Bryan et al., 20031.

[5] The scope of this study is to elucidate statistical features of convection in a global model and their resolution denendence, by conducting a series of high-resolution global simulations with resolution as fine as subkilometer that is finer regions through meridional circulations. Furthermore, convec-

- We are on the way to Zettascale computing. However, can Zettascale be the solution?
- it is not enough Simple aircraift in cruise would take billions of year!!

### Fundamental Questions:

• Is it possible to identify algorithms for PDE solutions that QC can perform effectively?



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- **2** Can a current QC speed up the solution of these multi-scale nonlinear partial differential equations, such as turbulent Navier-Stokes Equations (NSE)?



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- Is it currently possible with *noisy intermediate-scale quantum (NISQ)* architecture?



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- Is it currently possible with noisy intermediate-scale quantum (NISQ) architecture?
- 4 Can we really understand fundamental science with a QA for solving PDE?



• Navier-Stokes equations (NSE)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0; \tag{2}$$

$$\frac{\partial (\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) = -\nabla p + \nabla \cdot \mathbf{\emptyset} + \rho \mathbf{g}; \tag{3}$$

### Remarks

Continuum Hypothesis



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- Numerical Solution!!!

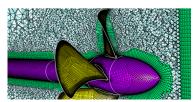


### Computational Fluid Dynamics

### What is it?

- Simulation of phenomena involving fluid flow;
  - 1 Domain representation;

### **Numerical Solution**



https://www.pointwise.com/

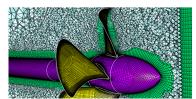


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https://www.pointwise.com/



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https://www.pointwise.com/

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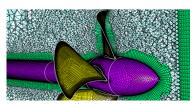
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- CFD simulations can be expensive;



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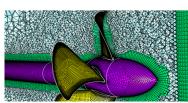
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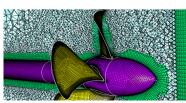
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- CFD simulations can be expensive;
- To test, design and analysis, several simulations can be needed;
- How to reduce computational costs, without disobeying the physics of problem?
- How Machine Learning can help?



## Neural Network: Brief description

### Machine Learning and Computer science

- AI: Looking for human-level intelligence;
- ML: Learn parttens in a dataset and performs predictions;
- Data Science: Find insights from data;

"To make progress in ML in science it is crucial that we incorporate computational models into the learning step."

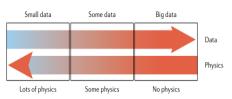
Prof. Amir Gholaminejad



### How to introduce physics into a neural network:

 Method 1; Where big data is available, but the governing physical law may not be known;

<sup>&</sup>lt;sup>a</sup>Karniadakis, G.E., Kevrekidis, I.G., Lu, L. et al. Nat Rev Phys 3, 422–440 (2021)

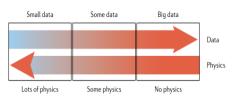




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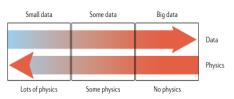




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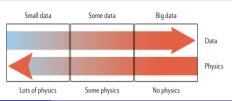




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- Method 4; where the physics is partially known and several scattered measurements are available (inverse problems) <sup>a</sup>

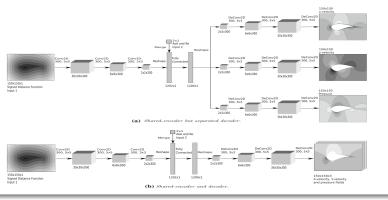
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# Physics-informed by a large number of data

### Model reduction by Machine Learning, Bhatnagar et 2019

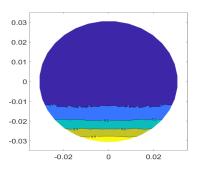


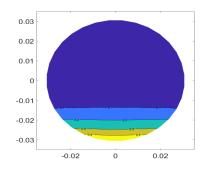
### Model reduction by Machine Learning

- Order reduction for aerodynamics;
- Convolutional neural network (ConvNet) with AutoEncoder;

# Physics-informed by a large number of data

Model reduction by Machine Learning, Ligang Lu and Kuochen Tsai (SHELL Oil Company), 2019.





### Model reduction by Machine Learning

- CFD model and predictions of oil water separation in horizontal pipelines;
- Convolutional neural network (CNN) model;

# Physics-informed by constraints

### Symmetries and Invariance

- Impose some properties to the trainning, such as: Symmetries and Invarience;
- Symmetry in a shear stress tensor;
- Invariance related to a reference or even conservation laws:

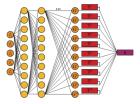


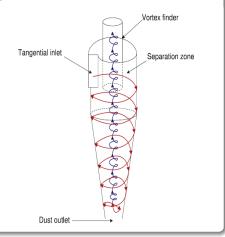
Fig. 2. Invariant, fully-connected (some connections are omitted for clarity), neural network architecture proposed by Ling et al. [17]. The circles indicate scalar values and the rectangles represent 3 x 3 second-order tensors.

Figure: Symetry and Invariance of reynolds stress tensor based on Invariance expansion; 1

<sup>&</sup>lt;sup>1</sup>Ling, J., Kurzawski, A., Templeton, J. (2016). Journal of Fluid Mechanics, 807, 155-1668 **○CE** doi:10.1017/jfm.2016.615

# Physics-informed by constraints

# Cyclone Separator, Caturwati Ni Ketut et al. 2017



### Number of samples

- 300 Simulations with different Re numbers;
- Data extracted from the cell centers of the mesh

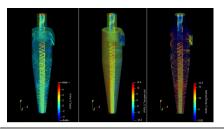
### Case Setup

- OpenFOAM (OF) simulation with SimpleFOAM solver;
- $\kappa \epsilon$  with curvature correction in OF, Spalart-Shur, 2000;

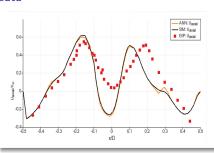


# Physics-informed by constraints

### Velocity streamlines



# Velocity profile versus experimental data



### Conservation Equation Table

	$R_{\mathbf{M},x}$	$R_{\mathbf{M},y}$	$R_{\mathbf{M},z}$	$ abla \cdot \mathbf{U}$
$\overline{V}$	$-4,85 \cdot 10^{-9}$	$-2,31\cdot 10^{-8}$	$1,65 \cdot 10^{-8}$	$-7, 2 \cdot 10^{-7}$
max(V)	$2, 3 \cdot 10^{-3}$	$4,9 \cdot 10^{-5}$	$1, 1 \cdot 10^{-3}$	0,047
min(V)	$-1 \cdot 10^{-3}$	$-4,10\cdot 10^{-5}$	$-1,8\cdot 10^{-3}$	-0,04

# Physics-informed by constraints

### Computational resource for simulation

Intel Core i7, 12 cores; 32GB RAM;

### Computational resource for PIDNN

- Intel Xeon® E5-2640 v4 2,4GHz
- Tesla P100 (Pascal), 3584 CUDA cores, 16GB VRAM;

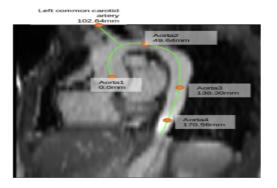
#### Computational Time

- **1** CFD Simulation computational Time  $\approx 1.6$  h
- 2 Training time  $\approx 3$  h
- **3** PIDNN computation  $\approx 1$  s
- **4** Computational "speedup"  $\approx 5000$ ;

<sup>&</sup>lt;sup>a</sup>L.H. Queiroz, F.P. Santos, J.P. Oliveira, M.B. Souza, Digital Chemical Engineering, Volume 1, 2021, 100002.

# Physics-informed Neural Network approximation

Flow through the aorta/carotid bifurcation of a healthy human subject: Positions of acquired 4D flow MRI measurements in the aorta/carotid bifurcation of a healthy volunteer.<sup>2</sup>

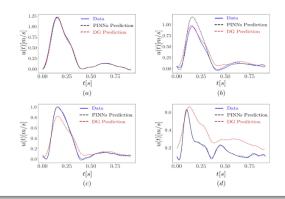


<sup>&</sup>lt;sup>2</sup>G. Kissas, Y. Yang, E. Hwuang et al. / Computer Methods in Applied Mechanics and Engineering 358 (2020) 112623 21



# Physics-informed Neural Network approximation

Comparison of the clinically acquired waveforms of blood velocity versus the predictions of the proposed physics-informed neural networks, and of a conventional Discontinuous Galerkin solver.<sup>3</sup>



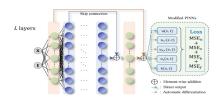
<sup>&</sup>lt;sup>3</sup>G. Kissas, Y. Yang, E. Hwuang et al. / Computer Methods in Applied Mechanics and Engineering 358 (2020) 112623 21



# Physics-informed Neural Network approximation (inverse problem)

A physics-informed deep learning method for solving direct and inverse heat conduction problems of materials.<sup>4</sup>

Inverse problem Setup



$$\rho(u)Cp(u)\frac{\partial u}{\partial t} - k(u)\frac{\partial^2 u}{\partial x^2} = 0 \quad in \quad \Omega, t \in [0, t_{max}]$$
(4)

<sup>&</sup>lt;sup>4</sup>Zhili He, Futao Ni, Weiguo Wang, Jian Zhang. Materials Today Communications, Volumescri 28, 2021, 102719.

# Neural Network: Brief description

#### Neural Network

• Recognized as fundamental nonlinear function approximators. It can be represented as:  $R(\mathbf{w}, b) = \int L[\mathbf{y}, \phi(\mathbf{w}, \mathbf{y}, \mathbf{x}, b)] p(\mathbf{y}, \mathbf{x}) d\mathbf{x} d\mathbf{y}$ 

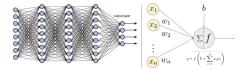


Figure: Deep neural network example

Deep Neural Network is actually a composition function!



# Neural Network: Brief description

#### Supervised neural network

- Depends on what the focus is on....
- We want to find a neural net,  $\phi(\mathbf{y}, \mathbf{x}; \theta(\mathbf{w}, b))$ , that maps  $\mathbf{x}$  to  $\mathbf{y}$
- Minimizing  $R(\mathbf{w}, b)$ , where the data are equiprobable;

$$R(\mathbf{w}, b) = \frac{1}{N} \sum_{i=0}^{N} \left[ y_i - \phi(y_i, \mathbf{x}, \mathbf{w}, b) \right]^2$$
(5)

being N the size of trainning data;

How we usually solve? Stochastic gradient decent

$$\theta_{k+1} = \theta_k - \alpha_k \frac{1}{m} \sum_{1}^{m} \nabla_{\theta} \left[ y_j - \phi(y_j, \mathbf{x}, \mathbf{w}, b) \right]^2$$
 (6)



# Deep Neural Network: Brief description

Let  $N^L(\mathbf{x}): \Re^{d_{in}} \to \Re^{d_{out}}$ . Let us say the weight matrix of the  $l^{th}$  layer and the bias vector are  $\mathbf{W}^l \in \Re^{N_l \times N_{l-1}}$  and  $\mathbf{b}^l \in \Re^{N_l}$ , respectively. Then, one can build the deep neural network below:

• 
$$N^0(\mathbf{x}) = \mathbf{x} \in \Re^{d_{in}}$$

where  $\phi$  is a composition function of  $N^{j}(\mathbf{x}), \quad j = 0 \cdots L$ 

$$\mathbf{u}(\mathbf{x}) \approx \mathbf{u}(\mathbf{x}; \theta) = \phi(\theta; \mathbf{x})$$
 (7)

where  $\theta = [\mathbf{W}, \mathbf{b}]$ 



# Deep Neural Network: Brief description

Let  $N^L(\mathbf{x}): \Re^{d_{in}} \to \Re^{d_{out}}$ . Let us say the weight matrix of the  $l^{th}$  layer and the bias vector are  $\mathbf{W}^l \in \Re^{N_l \times N_{l-1}}$  and  $\mathbf{b}^l \in \Re^{N_l}$ , respectively. Then, one can build the deep neural network below:

- $N^0(\mathbf{x}) = \mathbf{x} \in \Re^{d_{in}}$
- $N^l(\mathbf{x}) = \sigma\left(\mathbf{W}^l N^{l-1}(\mathbf{x}) + \mathbf{b}^l\right) \in \Re^{N_l}, \quad 1 < l \le L-1$

where  $\phi$  is a composition function of  $N^{j}(\mathbf{x}), \quad j = 0 \cdots L$ 

$$\mathbf{u}(\mathbf{x}) \approx \mathbf{u}(\mathbf{x}; \theta) = \phi(\theta; \mathbf{x})$$
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#### How to calculate derivatives?

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#### How to calculate derivatives?

- hand-coded analytical derivative;
- finite difference or other numerical approximations;
- symbolic differentiation (used in software programs such as Mathematica, Maxima, and Maple); and
- Automatic differentiation (the derivatives are evaluated using backpropagation);



#### How to calculate derivatives?

- a compositional function, then AD applies the chain rule repeatedly to compute the derivatives;
- There are two steps in AD: one forward pass to compute the values of all variables, and one backward pass to compute the derivative

Consider a FNN with one hidden layer two inputs  $(x_1, x_2)$  and one output (y);

$$v = -2x_1 + 3x_2 + 0.5 (8)$$

$$h = tahn(v) (9)$$

$$y = 2h - 1 \tag{10}$$



To calculate the derivative of  $\frac{\partial y}{\partial x_1}$  and  $\frac{\partial y}{\partial x_2}$  at (2,1):

Backward pass
$\frac{\partial y}{\partial u} = 1$
- 0
$\frac{\partial y}{\partial h} = \frac{\partial (2h-1)}{\partial h} = 2$
$\frac{\partial y}{\partial v} = \frac{\partial y}{\partial h} \frac{\partial h}{\partial v} = \frac{\partial y}{\partial h} \operatorname{sech}^2(v) \approx 1.573$
$\frac{\partial y}{\partial x_1} = \frac{\partial y}{\partial v} \frac{\partial v}{\partial x_1} = \frac{\partial y}{\partial v} \times (-2) = -3.146$
$\frac{\partial y_1^1}{\partial x_2} = \frac{\partial y}{\partial v} \frac{\partial v_1^1}{\partial x_2} = \frac{\partial y}{\partial v} \times 3 = 4.719$

Figure: AD calculation example

#### How to calculate derivatives?

- AD has two passes one forward pass and one backward:
- Any discretization derivative method has at least;
- It can be used for *n* order derivative <sup>a</sup>;

<sup>&</sup>lt;sup>a</sup>Lu, Lu and Meng, Xuhui and Mao, Zhiping and Karniadakis, George Em. SIAM Review, Volume 63,number 1, 208-228, 2021

# Universal approximation theorem

Approximation:  $Loss \rightarrow 0^{5}$ ;?

Let  $\sigma$  be any continuous sigmoidal function. The finite sum of the form:

$$G(x) = \sum_{j=1}^{N} \alpha_j \sigma(w_j x + b_j)$$
(11)

are dense in  $C(I_d)$ 

### Theorem (Pinks, 1999)

Let  $\mathbf{m}^i \in \mathbb{Z}_+^d$ ,  $i=1,\cdots,s$ , and set  $m=\max_{i=1,\cdots,s}|\mathbf{m}^i|$ . Assume  $\sigma \in C^m(\mathbb{R})$  and also is not a polynomial. Then the space of single hidden layer neural network:

$$\mathcal{M}(\sigma) = \sigma \left( \mathbf{W} \mathbf{x} + b \right) \tag{12}$$

is dense in  $C^{\mathbf{m}^1,\cdots,\mathbf{m}^s}(\mathbb{R}^d):=\bigcap_{i=1}^s C^i(\mathbb{R}^d).$ 

<sup>&</sup>lt;sup>5</sup>Kurt Hornik, Maxwell Stinchcombe, Halbert White, Multilayer feedforward networks are LAB CFD universal approximators, Neural Networks, Volume 2, Issue 5, 1989, Pages 359-366.

# Universal approximation theorem

### Universal approximation theorem

For  $\mathbf{m}=(m_1,\cdots,m_d)\in\mathbb{Z}^d_+$ , we set  $|\mathbf{m}|\!:=m_1+\cdots+m_d$  and

$$D^{\mathbf{m}} := \frac{\partial^{|\mathbf{m}|}}{\partial^{m_1} x_d \cdots \partial^{m_d} x_d} \tag{13}$$

We say  $f \in C^{\mathbf{m}}(\mathbb{R}^d)$  if  $D^{\mathbf{k}} f \in C^{\mathbf{d}} \ \forall \ \mathbf{k} \leq \mathbf{m}, \ \mathbf{k} \in \mathbb{Z}_+^d$ , where  $C^{\mathbf{d}} = \{f : \mathbb{R}^d \to \mathbb{R}\}$  is a space of continuous functions. We can recall the Pinkus' Theorem and say: For a any  $f \in C^{\mathbf{m}}(\mathbb{R}^d)$ , any compact  $K \subset \mathbb{R}^d$ , and any  $\epsilon > 0$ , there exists a  $g \in \mathcal{M}(\sigma)$  satisfying:

$$max|D^{\mathbf{k}}f - D^{\mathbf{k}}g| < \epsilon, \quad x \in K,$$
 (14)



# Universal approximation theorem

#### Universal approximation theorem

- Each neuron can be seen as a basis function;
- MLP with non-linear activations are universal functions approximators<sup>a</sup>;
- However, it says that MLP can approximate but...

<sup>a</sup>Kurt Hornik, Maxwell Stinchcombe, Halbert White, Multilayer feedforward networks are universal approximators, Neural Networks, Volume 2, Issue 5, 1989, Pages 359-366.

#### Limitation

Trainability;

**Questions:** Does the sequence of neural networks converge to the solution to PDFs?



### On the convergence of physics informed neural networks for linear second-order elliptic and parabolic type PDEs

Yeonjong Shin<sup>1</sup>, Jérôme Darbon<sup>1</sup>, and George Em Karniadakis<sup>1</sup>

Abstract. Physics informed neural networks (PINNs) are deep learning based techniques for solving partial differential equations (PDEs) encounted in computational science and engineering. Guided by data and physical laws, PINNs find a neural network that approximates the solution to a system of PDEs. Such a neural network is obtained by minimizing a loss function in which any prior knowledge of PDEs and data are encoded. Despite its remarkable empirical success in one, two or three dimensional problems, there is little theoretical justification for PINNs.

As the number of data grows, PINNs generate a sequence of minimizers which correspond to a sequence of neural networks. We want to answer the question: Does the sequence of minimizers converge to the solution to the PDE? We consider two classes of PDEs: linear second-order elliptic and parabolic. By adapting the Schauder approach and the maximum principle, we show that the sequence of minimizers strongly converges to the PDE solution in  $C^0$ . Furthermore, we show that if each minimizer satisfies the initial/boundary conditions, the convergence mode becomes  $H^1$ . Computational examples are provided to illustrate our theoretical findings. To the best of our knowledge, this is the first theoretical work that shows the consistency of PINNs.

AMS subject classifications: 65M12, 41A46, 35J25, 35K20

Key words: Physics Informed Neural Networks, Convergence, Hölder Regularization, Elliptic and Parabolic PDEs, Schauder approach

<sup>&</sup>lt;sup>1</sup> Division of Applied Mathematics, Brown University, Providence, RI 02912, USA

# Simple PINN-PDE example

#### Mathematical Formulation<sup>a</sup>

<sup>a</sup>I. Lagaris, A. Likas, D. Fotiadis. Artificial neural networks for solving ordinary and partial differential, 1997.Mathematics, Physics, Computer Science, Medicine IEEE transactions on neural networks

$$\frac{\partial g}{\partial x} = h(x, g), \quad g(0) = A$$
 (15)

With NN representation;

$$g(x) \approx g(x;\theta) = A + xNN(x;\theta);$$
 (16)

$$L(\theta) = \int_{\Omega_X} \left[ \frac{\partial g(x;\theta)}{\partial x} - h(x,g(x;\theta)) \right]^2 dx \tag{17}$$



# Simple PINN-PDE solution

• In a Riemman point of view, we have:

$$L(\theta) = \lim_{N \to \infty} \sum_{i=0}^{N} \left[ \frac{\partial g(x_i; \theta)}{\partial x} - h(x_i, g(x_i; \theta)) \right]^2$$
 (18)

Solving the problem with Stochastic gradient descent

$$\theta_{k+1} = \theta_k - \alpha_k \frac{1}{N} \sum_{i=1}^{N} \nabla_{\theta} \left[ \frac{\partial g(x_i; \theta)}{\partial x} - h(x_i, g(x_i; \theta)) \right]^2$$
 (19)

Finally,

$$g(x;\theta) = A + xNN(x;\theta); \tag{20}$$



# Forward PINN-PDE

#### Generalization PINN-PDE

Extension for General PDEs<sup>a</sup>

<sup>a</sup>M Raissi, P Perdikaris, GE Karniadakis Journal of Computational Physics 378, 686-707, 2017

#### PDE general form

Consider the general system PDE as:

$$\mathbf{f}(\mathbf{x}; \frac{\partial u_1}{\partial x_1}, \dots, \frac{\partial^2 u_m}{\partial x_1 \partial x_d}, \dots; \lambda) = 0, \quad \mathbf{x} \in \Omega,$$
(21)

with boundary conditions:

$$B(\mathbf{x}, \mathbf{u}) = 0$$
 on  $\partial \Omega$  (22)



# Forward Generalization PINN-PDE

#### Generalization PINN-PDE

$$\Im(\theta, \tau) = \omega_f \Im_f(\theta, \tau_f) + \omega_b \Im_b(\theta, \tau_b)$$
(23)

where,

$$\Im_f(\theta, \tau_f) = \frac{1}{|\tau_f|} \sum_{\mathbf{x} \in \tau_f} ||\mathbf{f}(\mathbf{x}; \frac{\partial \hat{u}_1}{\partial x_1}, \dots, \frac{\partial^2 \hat{u}_m}{\partial x_1 \partial x_d}, \dots; \lambda)||_2^2$$
 (24)

and

$$\Im_b(\theta, \tau_b) = \frac{1}{|\tau_b|} \sum_{\mathbf{x} \in \mathcal{T}} ||B(\mathbf{x}, \mathbf{u})||_2^2$$
 (25)



# Forward Generalization PINN-PDE

#### Generalization PINN-PDE

PINN algorithm<sup>a</sup>;

#### **Procedure 2.1** The PINN algorithm for solving differential equations.

- Step 1 Construct a neural network  $\hat{u}(\mathbf{x}; \boldsymbol{\theta})$  with parameters  $\boldsymbol{\theta}$ .
- Step 2 Specify the two training sets  $\mathcal{T}_f$  and  $\mathcal{T}_b$  for the equation and boundary/initial conditions.
- Step 3 Specify a loss function by summing the weighted  $L^2$  norm of both the PDE equation and boundary condition residuals.
- Step 4 Train the neural network to find the best parameters  $\theta^*$  by minimizing the loss function  $\mathcal{L}(\theta; \mathcal{T})$ .

Figure: Physics-informed machine learning for PDE solution;



<sup>&</sup>lt;sup>a</sup>Lu, Lu and Meng, Xuhui and Mao, Zhiping and Karniadakis, George Em. DeepXDE: A deep learning library for solving differential equations.SIAM Review volume 63, number 1, 208-228, 2021

### States: Finite Hilbert space representation....

- States are vectors;
- Operations are matrices;



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- **Qubit** Superposition  $|\psi\rangle = \alpha_0|0\rangle + \alpha_1|1\rangle = \begin{pmatrix} \alpha_0 \\ \alpha_1 \end{pmatrix} \in \mathbb{C}^2$  (ket representation) and  $\langle \psi |$  dual Hilbert Space (bra representation)



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- 3  $\langle \psi | \phi \rangle$  inner product and  $\rangle \psi | \phi \langle$  outer product, where  $| \phi \rangle = \beta_0 | 0 \rangle + \beta_1 | 1 \rangle$



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- **6** The quantum state  $|\psi\rangle$ , one measures the probability of the state  $|k\rangle$  as  $|\alpha\rangle$

# Quantum State Operations

### Operations: Logic Gates

#### Circuit

- Classical Boolean circuit (AND, OR and NOT gates on an n-bit.
- Quantum circuit uses unitary quantum gates
- NOT-gates

$$NOT = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Madamard

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix}$$

3 Interference:

$$\frac{1}{\sqrt{2}}\begin{pmatrix}1 & -1\\1 & 1\end{pmatrix}\frac{1}{\sqrt{2}}\left(|0\rangle + |1\rangle\right) = \frac{1}{2}\left(|0\rangle + |1\rangle - |0\rangle + |1\rangle\right)$$



# Quantum Circuit

### **Bloch Sphere:** Graph representation

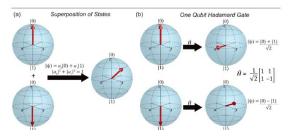


Figure: Graph representation<sup>6</sup>

$$|\psi\rangle = \cos(\theta/2)|0\rangle - e^{i\phi}\sin(\theta/2)|1\rangle$$
 (26)

where  $0 < \theta \le 2\pi$  and  $0 < \phi \le \pi$ 

<sup>6</sup>Moreno-Pineda, Eufemio et al. Molecular spin qudits for quantum algorithms Chemical [ДВ⊖СГП Society Reviews, 47. 2018.

# Quantum Circuit

### Quantum Circuit: A set of Gates operations

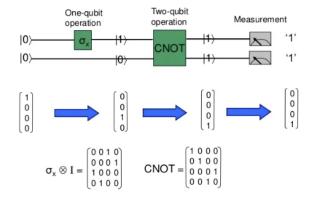


Figure: Quantum Circuit Example<sup>7</sup>

<sup>&</sup>lt;sup>7</sup>Moreno-Pineda, Eufemio et al. Molecular spin qudits for quantum algorithms Chemical [AB**SCFI**] Society Reviews, 47, 2018.

# Quantum computer steps

f 0 Start with all qubits with preparable state (init with all  $|0\rangle$ )

### Quantum parallelism

$$U\left(\frac{1}{2^{n/2}}\sum|x\rangle|0\rangle\right) = \frac{1}{\sqrt{2^n}}\sum|x\rangle|f(x)\rangle \tag{27}$$

- This has all 2<sup>m</sup> function values!
- However it produces one random  $|x\rangle|f(x)\rangle$ : Means that all will be lost!



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• Suppose a classical computation  $f:[0,1]^n \to [0,1]^m$ 

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#### Quantum parallelism

- Suppose a classical computation  $f:[0,1]^n \to [0,1]^m$
- Convert this to quantum circuit  $U:|x\rangle|0\rangle \rightarrow |x\rangle|f(x)\rangle$

$$U\left(\frac{1}{2^{n/2}}\sum|x\rangle|0\rangle\right) = \frac{1}{\sqrt{2^n}}\sum|x\rangle|f(x)\rangle \tag{27}$$

- This has all  $2^m$  function values!
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### General discretization by finite volume

Let's say:

$$\mathbf{U} = [\rho, \rho v_1, \rho v_2, \rho v_3, \epsilon_T]^T, \quad \mathbf{F} = [\rho v_1, \tau_{1,1}, \tau_{1,2}, \tau_{1,1}, \epsilon_1]^T$$

$$\mathbf{B} = [\rho v_2, \tau_{2,1}, \tau_{2,2}, \tau_{2,1}, \epsilon_2]^T, \quad \mathbf{C} = [\rho v_3, \tau_{3,1}, \tau_{3,2}, \tau_{3,1}, \epsilon_3]^T, \quad \mathbf{J} = \mathbf{0}$$

In this sense, the N-S equations can be writen as:

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x_1} + \frac{\partial \mathbf{B}}{\partial x_2} + \frac{\partial \mathbf{C}}{\partial x_3} = \mathbf{J}$$
 (28)



## General discretization by finite volume

#### For example:

$$\frac{\partial \mathbf{F}}{\partial x_1} \approx \frac{\mathbf{F}_{I+1} - \mathbf{F}_I}{\Delta x_1} \tag{29}$$

- one billion of nodes! Means billions degrees of freedom (At least)
- Lets says around 30 qubits!!

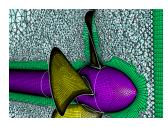


Figure: Volume discretization (mesh) https://www.pointwise.com/

It can lead to linear system of equation:

$$\mathbf{A}\mathbf{y} = \mathbf{b} \tag{30}$$

Most of the quantum algorithm are based on the solution of this linear system of equations.

### Current work applied to fluid dynamics

npj Quantum Information

www.nature.com/npiol

ARTICLE

Finding flows of a Navier-Stokes fluid through quantum computing

Frank Galtan (2) 15

There is great interest in using quantum computers to efficiently simulate a quantum system's dynamics as existing classical computers cannot do this. Little attention, however, has been given to quantum simulation of a classical nonlinear continuum system such as a viscous fluid even though this too is hard for classical computers. Such fluids obey the Navier-Stokes nonlinear partial differential equations, whose solution is essential to the aerospace industry, weather forecasting, plasma magnetohydrodynamics, and astrophysics, Here we present a quantum algorithm for solving the Navier-Stokes equations. We test the algorithm by using it to find the steady-state inviscid, compressible flow through a convergent-divergent nozzle when a shockwave is (is not) present. We find excellent agreement between numerical simulation results and the exact solution, including shockwave capture when present. Finally, we compare the algorithm's computational cost to deterministic and random classical algorithms and show that a significant speed-up is possible. Our work points to a large new application area for quantum computing with substantial economic impact, including the trillion-dollar aerospace industry, weather-forecasting, and engineered-plasma technologies.

In one of the earliest papers motivating the development of quantum computers Feynman' pointed out that such computers are better suited to simulate the dynamics of a quantum system than existing classical digital computers. This is because classical computers require computational resources that scale exponentially with the size of the simulated quantum system, while quantum computers do not. It is widely expected that quantum simulation of quantum systems will be a major application area for

npi Quantum Information (2020)6:61; https://doi.org/10.1038/s41534-020-00291-0

quantum computing. Quantum systems are not unique in being difficult to simulate. Many important classical systems exist whose dynamics strongly couple many degrees of freedom over multiple length-scales. making their simulation with classical computers hard. An example is a viscous fluid whose flows satisfy the Navier-Stokes nonlinear partial differential equations PDEs). Solving these PDEs is the primary task for such disease problems as serospace. flight vehicle design, weather-forecasting, probing the dynamics of turbulence, and determining the magneto-hydrodynamics of plasmas in space and in earth-bound technologies. In spite of its importance, little work has been done on finding quantum

thus extends to a general quantum algorithm for solving PDEs. To inviscid compressible flow through a convergent-divergent nozzle. Such nozzles are used in rocket engines and in jet engines. for supersonic flight and display a rich range of flows, from smooth acceleration from subsonic to supersonic speeds to the appearance of a normal shockwave in the divergent part of the nozzle, which causes supersonic flow to decelerate to subsonic speeds as it crosses the narrow shock region. We find excellent agreement between numerical simulation of our algorithm and the exact flow solution, including shockwave capture when present. We then compare the quantum algorithm's computational cost to deterministic and random classical algorithms, and discuss flow regimes where there is a computational advantage. Lastly, we point to the large new application area our quantum PDF algorithm opens up for quantum computing which includes many economically important problems. The quantum algorithm we present is independent of the previously cited papers, and of a quantum algorithm for elliptic PDEs

Contents lists available at ScienceDirect Computers and Fluids journal homepage: www.elsevier.com/locate/compfluid

Parallel evaluation of quantum algorithms for computational fluid dynamics\*

René Steiil14. George N. Barakos2

University of Chargow: School of Engineering, James West South Building, C12 8QQ, UR ABSTRACT

ARTICLE INFO Article history

work is limited to a single CFD algorithm. However, building on this work a broader cases of algorithms

The development and evaluation of quantum computing algorithms for computational fluid dynamics is described. A hybrid classical quantum hardware approach is assumed where selected computationally intensive parts of the solver are implemented as quantum circuits. The vortex-in-cell method is considered as an example where the Quantum Fourier Transform is used to build a Poisson solver. Computational aspects of simulating the required quantum circuits on a classical parallel computer are discussed includ-ing an analysis of the required data exchanges for a distributed-memory parallelization using messagepassing. The effect of errors and noise in the quantum algorithm on the flow solution is analyzed and it is shown that despite inevitable noise and uncertainties, meaningful flow simulations can be performed using a hybrid classical/quantum hardware approach. An improved version of the vortex-in-cell method with increased resilience to noise is also discussed along with suggestions for future steps. The presented

1. Introduction

In recent years the fleid of quantum computing (QC) [1] has grown into an active and diverse field of research and significant progress has been made with building quantum computers. For a small number of annications, quantum algorithms have been developed that would lead to a significant speed-up relative to classical methods when executed on a suitable quantum computer Despite this research effort, proceess in defining suitable applications for quantum computers has been relatively limited and two posite integers and Grover's algorithm for quantum search are still among the main applications. Further applications have been de-

outational fluid dynamics applications. In the absence of the required quantum hardware, targe-scate parattet simulations on parallel classical computers are required in developing such algorithms. A hybrid classical/ouantum hardware approach is assumed

Quantum computer applications aim to achieve a computational effects. The first and most important for the algorithms considered here is quantum parallelism, based on the use of oubits representing a superposition of infinitely many possible states of a binary other main quantum effects used to create computational canabil-

In this two case, the PDEs are discretized<sup>8</sup>. The system of equations produzed is solved either by HHL or a Variational Quantum algorithms<sup>9</sup>.

<sup>8</sup>René Steijl, George N. Barakos, Computers & Fluids, Volume 173, 2018, Pages 22-28.



Gaitan, F. npj Quantum Inf 6, 61 (2020)

### Current work applied to fluid dynamics

Towards Solving the Navier-Stokes Equation on Quantum Computers

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#### abstract

In this paper, we explore the suitability of upcoming novel computing technologies, in particular adiabatic annealing based quantum computers, to solve fluid dynamics problems that form a critical component of several science and engineering applications. We start with simple flows with well-studied flow properties, and provide a framework to convert such systems to a form amenable for deployment on such quantum annealers. We analyze the solutions obtained both qualitatively and quantitatively as well as the sensitivities of the various solution selection schemes on the obtained solution.

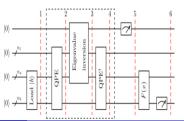
keywords: Quantum computing - quantum annealers - fluid dynamics - turbulence - linear systems.

# Adiabatic annealing-based quantum computers (quenching Ising type model)<sup>10</sup>

<sup>10</sup> N. Ray, T. Banerjee, B. Nadiga, S. Karra, "Towards Solving the Navier-Stokes Equation on Quantum Computers," arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers," arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers," arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.09033v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.0903v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.0903v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers, arXiv:1904.0903v1 [cs.NA] Towards Solving the Navier-Stokes Equation on Quantum Computers (cs.NA) Towards Solving the Navier-Stokes Equation on Quantum Computers (cs.NA) Towards Solving the Navier-Stokes Equation on Quantum Computers (cs.NA) Towards (cs.

### Harrow-Hassidim-Lloy Algorithm <sup>a</sup>

$$\mathbf{0} \ \mathbf{A} | \mathbf{y} \rangle = | \mathbf{b} \rangle$$

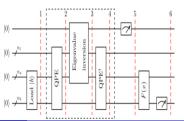




### Harrow-Hassidim-Lloy Algorithm <sup>a</sup>

$$\mathbf{0} \ \mathbf{A} | \mathbf{y} \rangle = | \mathbf{b} \rangle$$

$$\mathbf{2} \ \mathbf{A} = \sum_{j=0}^{N-1} \lambda_j |u_j\rangle \langle u_j|, \quad \lambda_j \in \Re$$



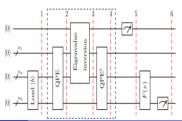


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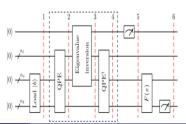
3 
$$\mathbf{U} = e^{i\mathbf{A}t} = \sum_{j=0}^{N-1} e^{i\lambda_j t} |u_j\rangle\langle u_j|$$





### Harrow-Hassidim-Lloy Algorithm <sup>a</sup>

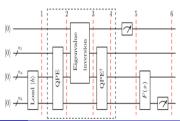
- $\mathbf{0} \ \mathbf{A} | \mathbf{y} \rangle = | \mathbf{b} \rangle$
- **2**  $\mathbf{A} = \sum_{j=0}^{N-1} \lambda_j |u_j\rangle\langle u_j|, \quad \lambda_j \in \Re$
- 3  $\mathbf{U} = e^{i\mathbf{A}t} = \sum_{j=0}^{N-1} e^{i\lambda_j t} |u_j\rangle\langle u_j|$
- **4**  $\mathbf{A}^{-1} = \sum_{j=0}^{N-1} \lambda_j^{-1} |u_j\rangle\langle u_j|$





### Harrow-Hassidim-Lloy Algorithm <sup>a</sup>

- $\mathbf{1} \mathbf{A} | \mathbf{y} \rangle = | \mathbf{b} \rangle$
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- **5**  $|\mathbf{y}\rangle = \mathbf{A}^{-1}|\mathbf{b}\rangle = \sum_{i=0}^{N-1} \lambda_i^{-1} b_i |u_i\rangle$

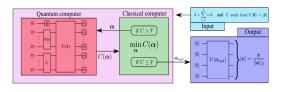




### Variational Quantum Linear Solver

### Variational Quantum Linear Solver (VQLS)

- Tested with problem size of 1024 × 1024.<sup>a</sup>;
- The input to VQLS is a matrix A written as a linear combination of unitaries  $\bf A$  and a short-depth quantum circuit U which prepares the state  $|{\bf b}\rangle$
- The hybrid quantum-classical optimization loop until the cost is below a user-specified threshold
- When the loop terminates, one obtains  $|\mathbf{y}\rangle = \mathbf{y}/||\mathbf{y}||_2$





<sup>&</sup>lt;sup>a</sup>Carlos Bravo-Prieto et al. 2020.

## Quantum Annealing Linear Solver

#### Quantum Annealing Linear Solver

- Transform in a binary matrix;
- Quadratic unconstrained binary optimization (QUBO) specific for DWave QPU

$$\frac{\partial u_x}{\partial t} = -\frac{1}{\rho} \frac{\partial \rho}{\partial x} + \mu \frac{\partial^2 u_x}{\partial x^2} \tag{31}$$

After discretization:

$$\mathbf{A}\mathbf{y} = \mathbf{b}, \quad y_j = \sum_{j=1}^{np} 2^{j_0 - j} q_j, \quad \mathbf{A}^q \mathbf{q} = \mathbf{b}$$
$$\bar{q} = min||\mathbf{A}^q \mathbf{q} - \mathbf{b}||_2^2$$
(32)

$$f(\mathbf{q}) = \sum_i \mu_i q_i + \sum_{i < j} \omega_{i,j} q_i q_j, \quad \mu_i = \sum_i A^q_{i,j} (A^q_{i,j} - 2b_i), \quad \omega_{i,k} = 2 \sum_i A^q_{i,j} A^q_{i,k} \mathcal{L}^q_{i,k} \mathcal{$$

### How about Neural network can help?

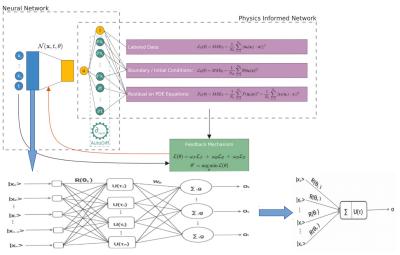
How about Neural network in quantum computer?

This raises another question: how Neural network can help?



## Methodologies

### Quantum Physics-informed Neural Networks Description<sup>11</sup>:





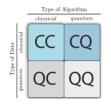


## Quantum Neural Network description

#### QNN and NN similarities

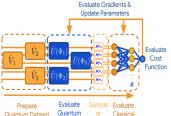
- Input data for training and testing
   Can be classical or quantum data<sup>a</sup>.
- Struture Input Layer, Hidden layers and output layers - Similar for QNN

ahttps://www.tensorflow.org/quantum



#### QNN and NN similarities

- Based on  $\theta$  **Similar for QNN**
- Non-linearity and optimization -Hybrid step - Quantum and classical
- Backpropagation Hybrid step -Quantum and classical



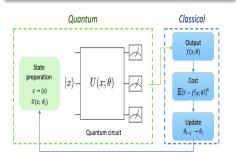
Model



## Quantum Neural Network description

#### Quantum neural network

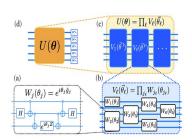
• General form  $QNN(x, \theta)$ :



#### Quantum circuit

• Quantum circuit representation<sup>a</sup>:

<sup>a</sup> Abbas, A., Sutter, D., Zoufal, C. et al. The power of quantum neura https://doi.org/10.1038/s43588-021-00084-1





## Quantum PINN formulation

#### PDE general form

Consider the general system PDE as:

$$\mathbf{f}(\mathbf{x}; \frac{\partial u_1}{\partial x_1}, \dots, \frac{\partial^2 u_m}{\partial x_1 \partial x_d}, \dots; \lambda) = 0, \quad \mathbf{x} \in \Omega,$$
(34)

with boundary conditions:

$$B(\mathbf{x}, \mathbf{u}) = 0 \quad on \quad \partial\Omega \tag{35}$$

$$I(\mathbf{x}, \mathbf{u}, \lambda), \quad for \quad \mathbf{x} \in \tau_j$$
 (36)

#### Can we formulate a Quantum PINN based on Quantum Circuit?

• Let's say:

$$u(\mathbf{x}) \approx QNN(\mathbf{x}; \theta)$$
 (37)

• We can rewrite the PDE general form as a fuction of the QC because we could compute:  $\frac{\partial QNN(\mathbf{x};\theta)}{\partial \mathbf{x}}$  - automatic shifted differentiation quantum



## Universal approximation for Quantum Neural Network

Working Hypothesis: QNN can always be used as an universal approximation.

#### Tools

- QNN libraries against their classical counterpart for regression:
  - TensorFlow Quantum (TFQ) based on Cirg.
  - IBM QNN based on Qiskit.
- Implementation of a specialized Quantum Neural Networks:
  - Quantum Convolutional Neural Networks (QCNN) with Bayesian Learning Initialization (BLI).
  - Avoid barren plateau (BP).

Measurement: (i) Memory capacity (ii) Generalization power (iii) Speed-up? (iv) Fisher information matrix.

**Expected**: A better understanding of QNN for regression and absence of barren plateau.

Risk Assessment: Medium-high risk

Significance: Overcome the barren plateaus.



## A Quantum Physics-Informed Neural Network (QPINN)

Working Hypothesis: QPNN can solve different type of non-linear PDE.

#### Challenges

- Can it be done by QNN libraries based: TFQ and QNNQ.
- Specialized Quantum PINN:
  - Quantum Convolutional Neural Networks (QCNN) with Bayesian Learning Initialization (BLI).
  - Avoid barren plateau (BP).



**Expected**: a PDE quantum solver. Risk Assessment: Medium-high risk

**Significance**: This might help to solve several currently impossible problems to answer that demand intensive computation.

### Global Risk

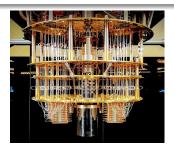
#### Global Risk

1 These algorithms depend on the development of a more powerful quantum processors.

### Noise Intermediate Scale Quantum (NISQ) era:

- Limitation of current quantum computers:
  - 1 Limitation of number of qubits;

 Google has already proved quantum supremacy on a quantum computer, and several countries, companies, and research centers are investigating this new scientific paradigm.



IBM Q quantum computer.



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### Noise Intermediate Scale Quantum (NISQ) era:

- Limitation of current quantum computers:
  - 1 Limitation of number of qubits;
  - Noise sources limit the number of operations:
  - 3 Problems of error correction:

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## Questions? I have several questions!<sup>12</sup>



